

Foreshock electron beams and electrostatic waves observed by Cluster

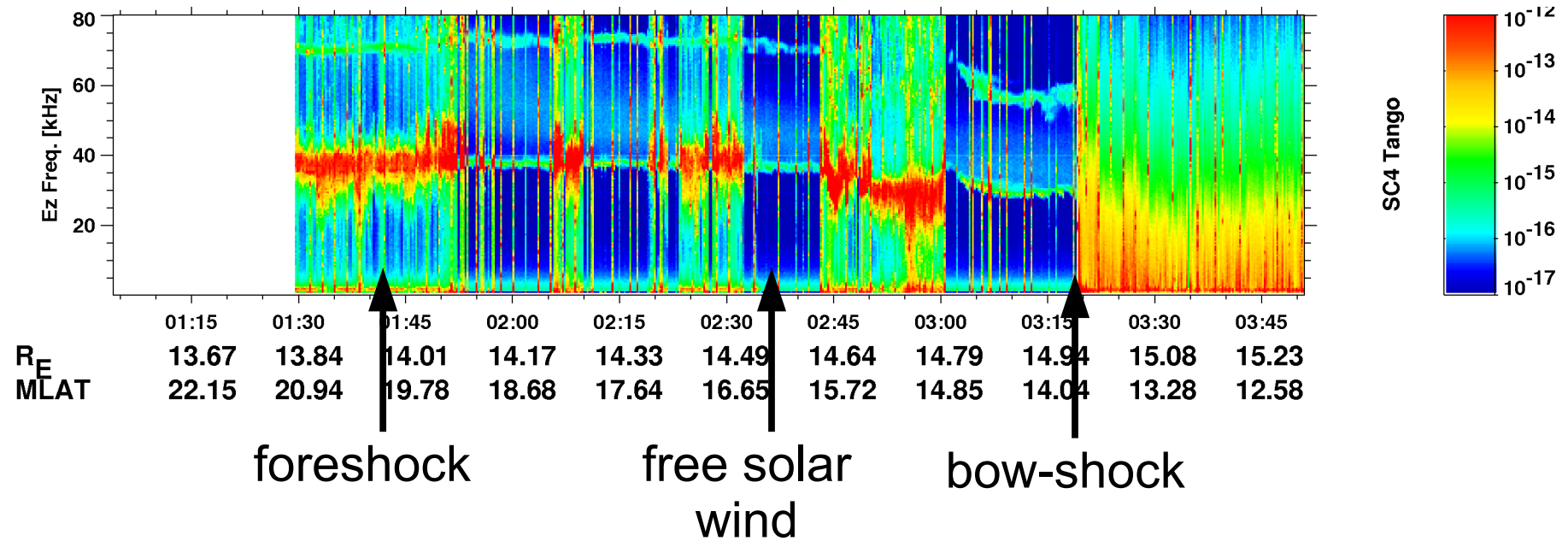
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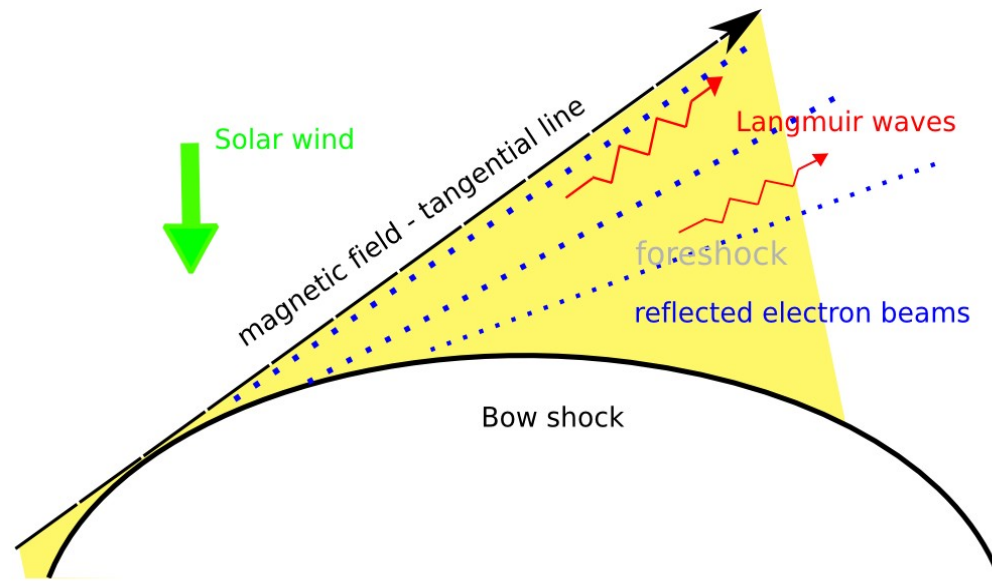
With thanks to Cluster WHISPER, FGM, CIS and PEACE instrument teams and to the Cluster-Active Science Archive.

Electron foreshock



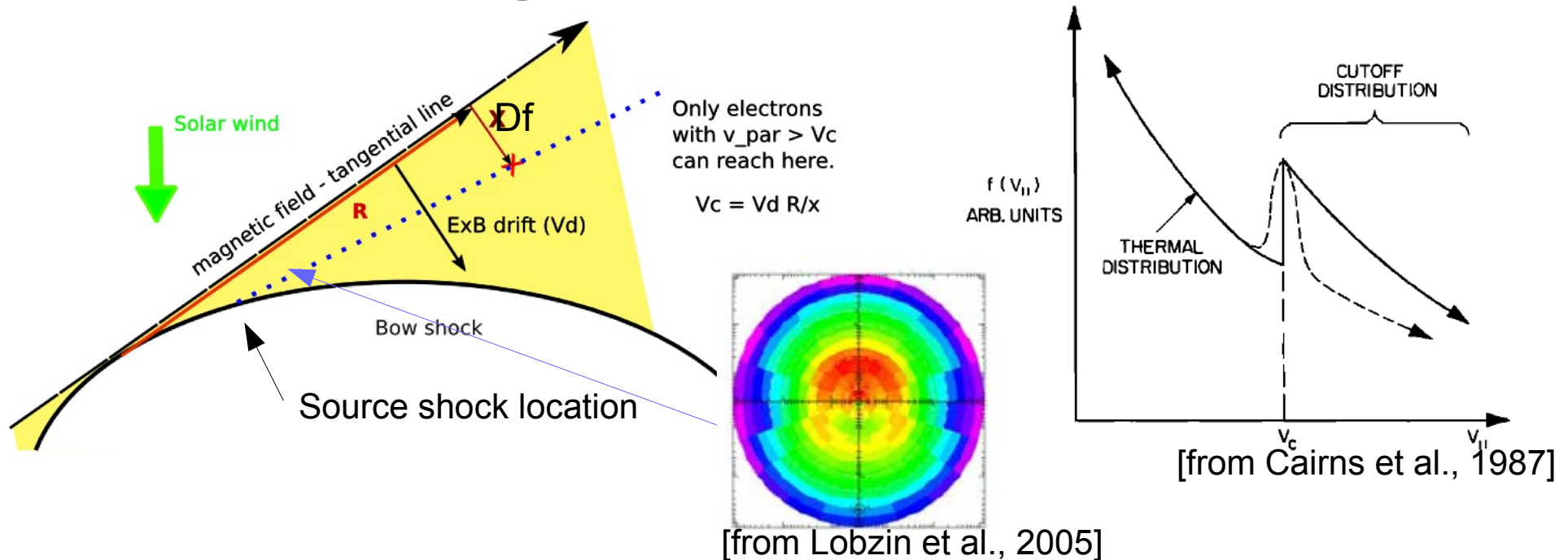
- Foreshock: a region of solar wind magnetically connected to the bow shock
- Populated by beams of back-streaming electrons (and further downstream ions), reflected and accelerated by the shock.
- Characteristic observable signature – intense electrostatic wave activity around the electron plasma frequency.
- Foreshock electron beams are a universal feature shocks: observed at planets, interplanetary shock, termination shock ...

Foreshock electron beams



- A fraction of incident solar wind electrons is reflected and accelerated by the bow shock [*Leroy & Mangeney, 1984, Wu, 1984*].
- These electrons (10 eV – 10's of keV) are responsible for the generation of the Langmuir waves via beam plasma instability.
- ExB drift in the foreshock acts as a velocity filter on the reflected electrons, forming the beams [*Fitzenreiter et al., 1990*].

Origin of the beams



Foreshock beam formation:

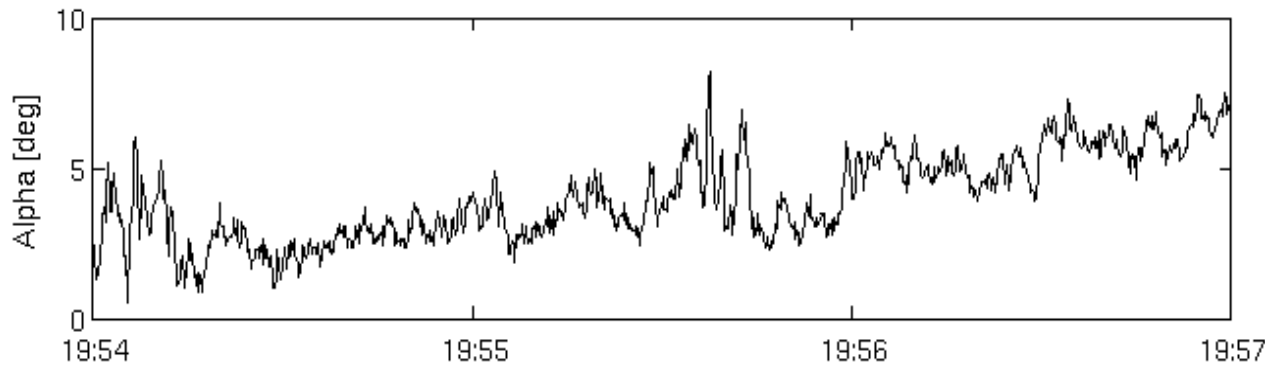
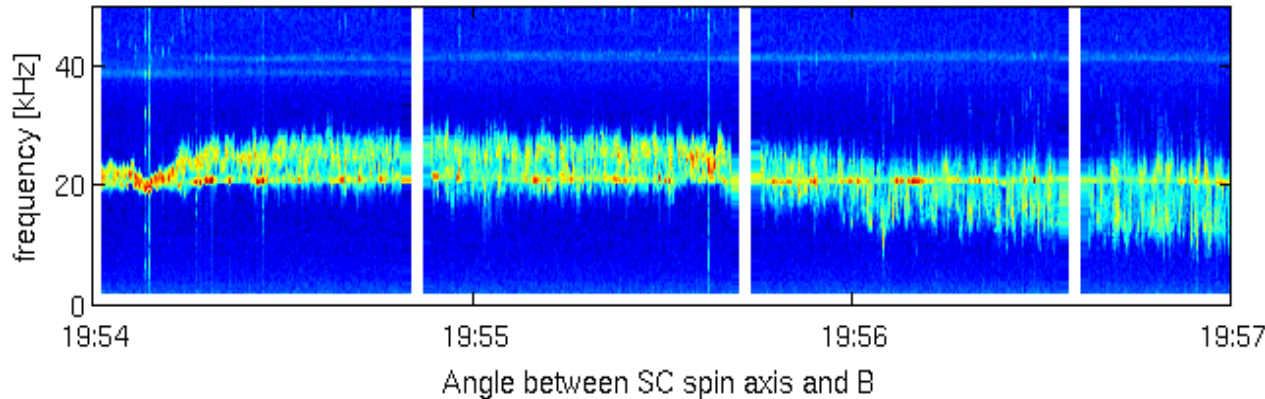
- Cutoff distribution [Filbert & Kellogg, 1979, Fitzenreiter et al., 1990]:
 - ➔ For each point (R, Df) cutoff velocity is defined $V_c = VE \times B R / Df$
 - ➔ “Bump on tail” required for the beam-plasma instability is formed at V_c
- Bump formed by reflection process [Leroy & Mangeney, 1984, Wu, 1984].
 - ➔ Beam with loss-cone (“rabbit ears”)
 - ➔ Model allows to calculate beam energy from parameters of source shock (θ_{Bn} , $B_{\text{down}}/B_{\text{up}}$ etc.)

Observation of foreshock electron beams

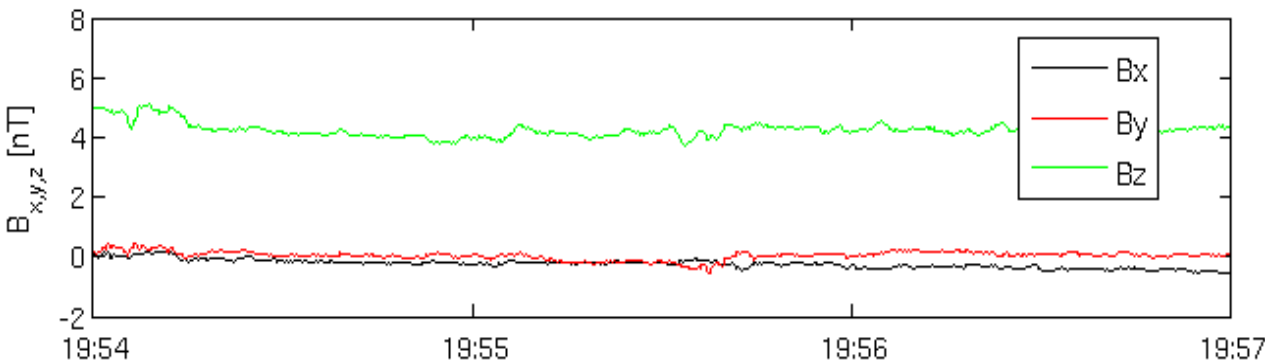
- Foreshock beams are typically faint ($n_b/n_e = 10^{-2} - 10^{-4}$) and non-stationary.
- Beams are present consistently, but not many direct observations were published [Fitzenreiter et al., 1996, Kasaba et al., 2000].
- Under favorable conditions, PEACE electron spectrometer on Cluster can be used to observe and characterize these beams - at least the slower ones.
 - Large geometrical factor needed (HEEA)
 - Good energy resolution (BM preferred)

Foreshock event C2 2011-04-21

WHISPER C2 2011-04-21 19:54:00.000



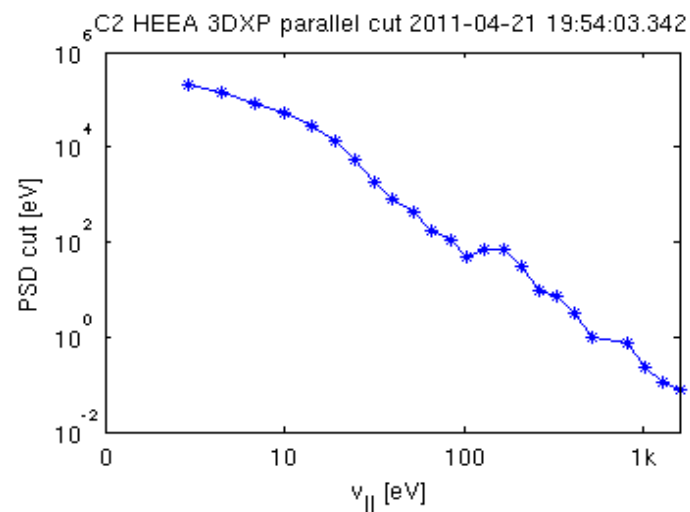
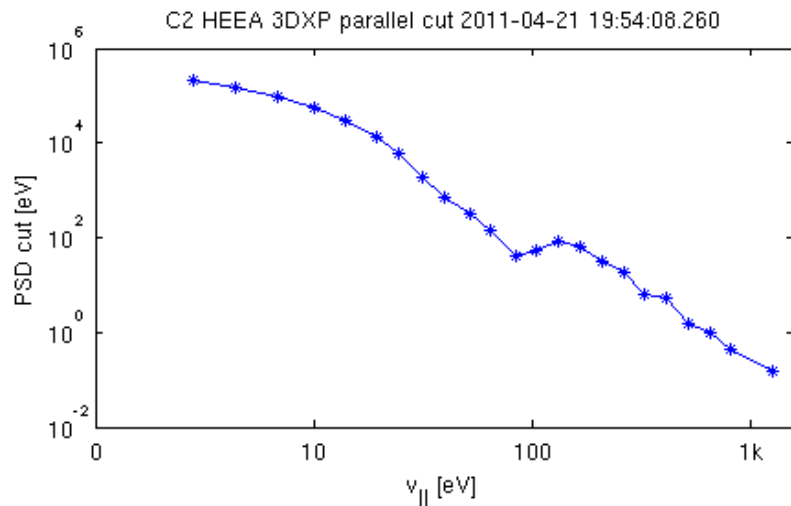
FGM B-field C2, tstart = 2011-04-21 19:54:00.000



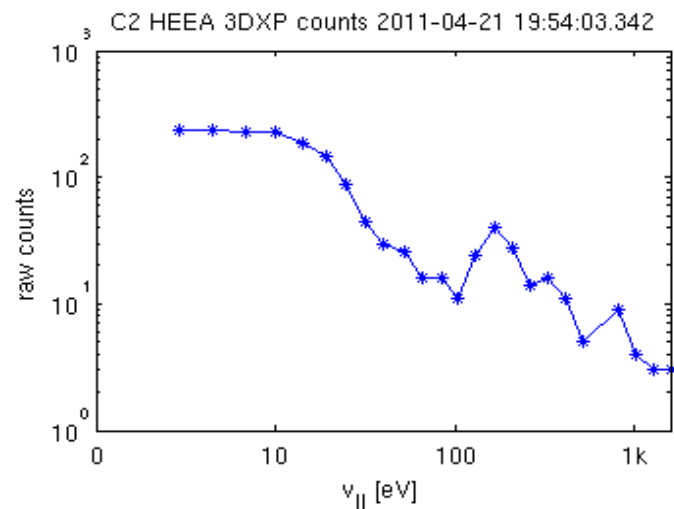
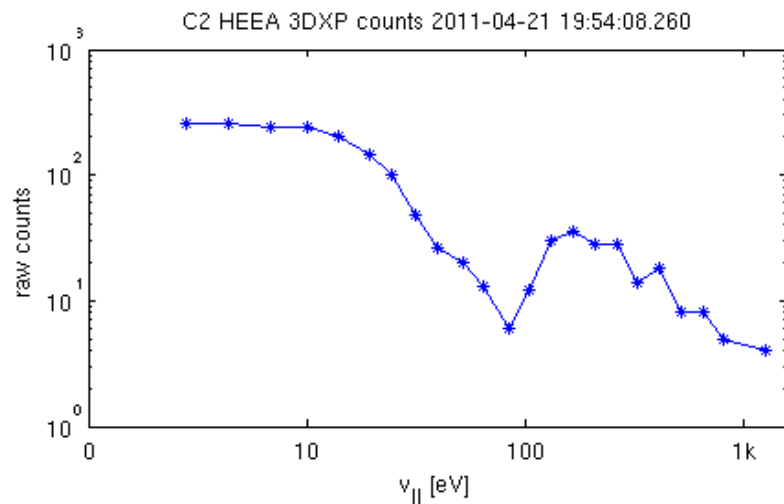
Suitable event:

- Cluster in BM
 - ➔ PEACE HEEA on
- B field in GSE Z direction (near parallel to spin axis)
 - ➔ pitch angle distribution available every PEACE sweep (8 per second).
- FS waves above and below ω_p observed

PEACE observed beams



B-parallel
PSD cut
(30 deg. bin)

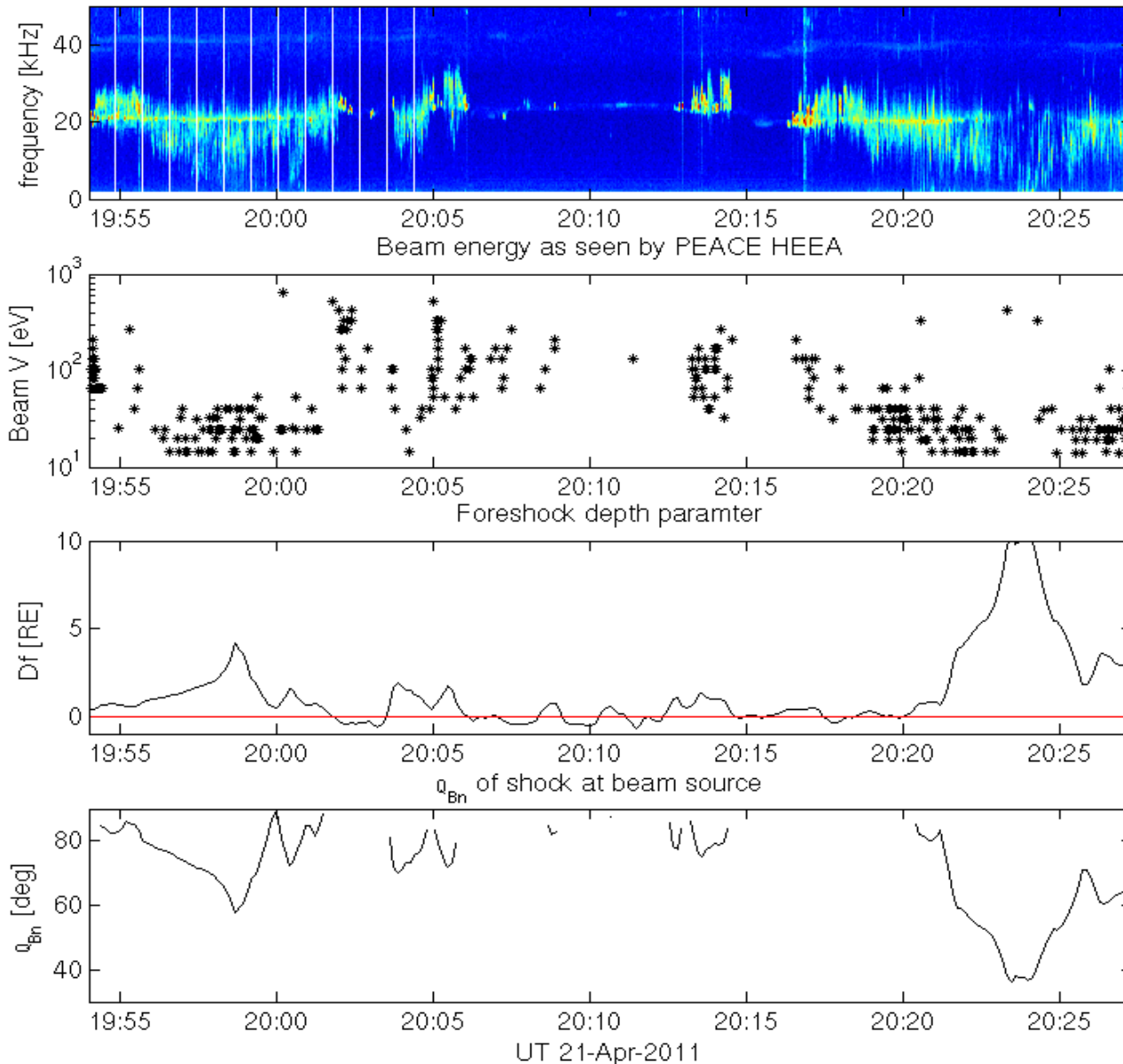


Raw
counts

- Beam-like bumps between 20 eV and 1 keV observed in the B-parallel bin
 - Higher energies not available due to low count rates
- ➔ In the ~ 100 eV range, the counting statistics is typically good.

Beam identification

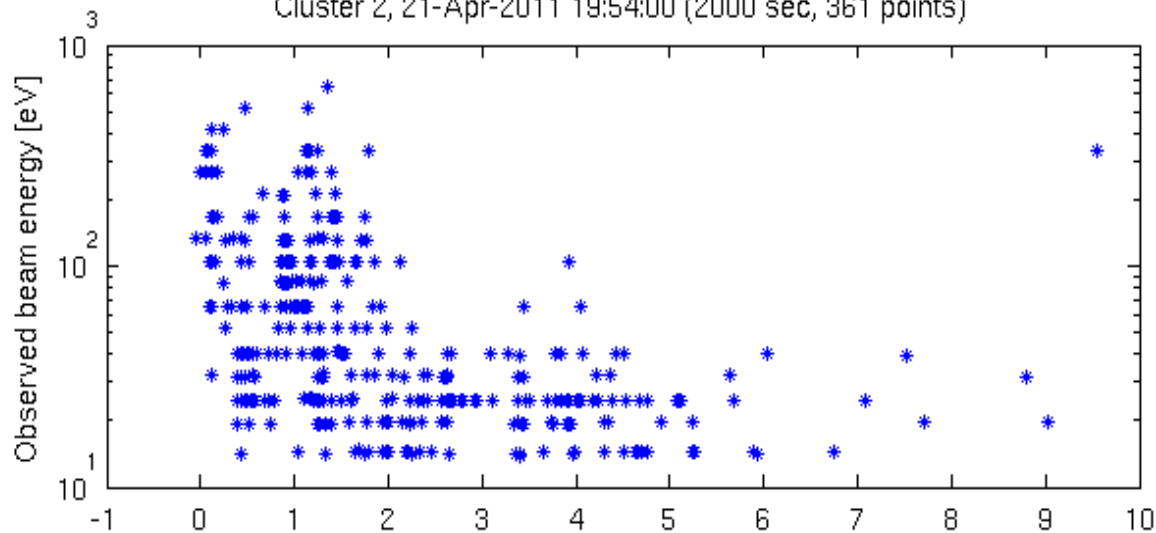
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- Beams identified automatically
 - ➔ Statistically significant ($> 2\sigma$) increase in PSD.
 - 361 beams between 13 eV and 850 eV identified ($T_e \sim 5$ eV).
 - Beam energy correlates with foreshock waves
 - ➔ Higher energy \rightarrow narrow-band
 - ➔ Low energy \rightarrow broad-band
- Consistent with [Lacombe et al., 1985, Cairns & Fung, 1988].
- Df and Θ_{Bn} parameters calculated.

Comparison to model - statistics

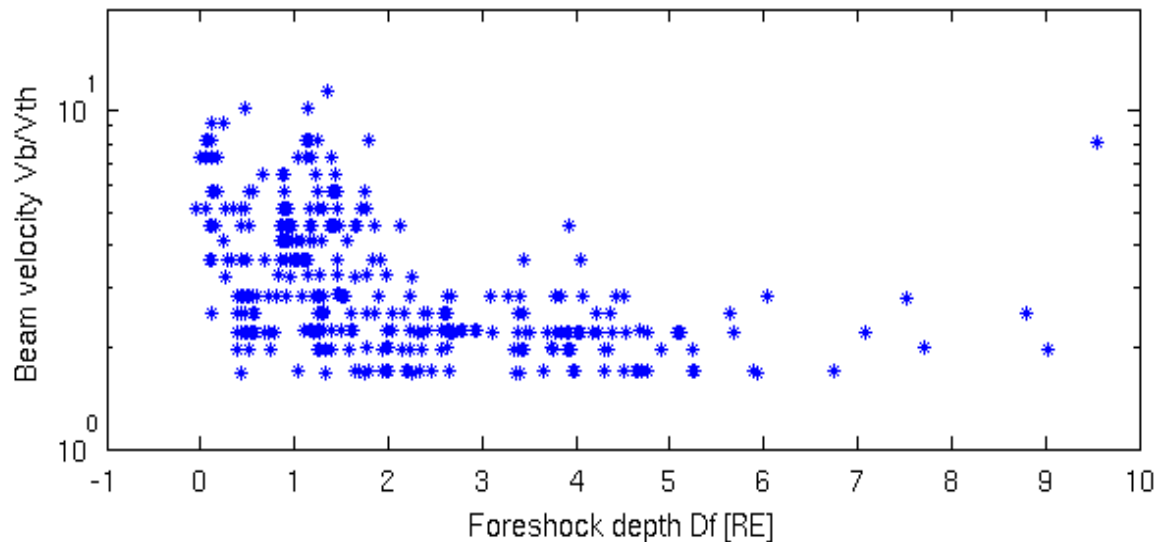
Cluster 2, 21-Apr-2011 19:54:00 (2000 sec, 361 points)



Strong dependence of beam energy on foreshock depth parameter [Fitzenreiter et al., 1990].

➔ Energetic beams $E > 100$ eV found for $D_f < 2$ RE.

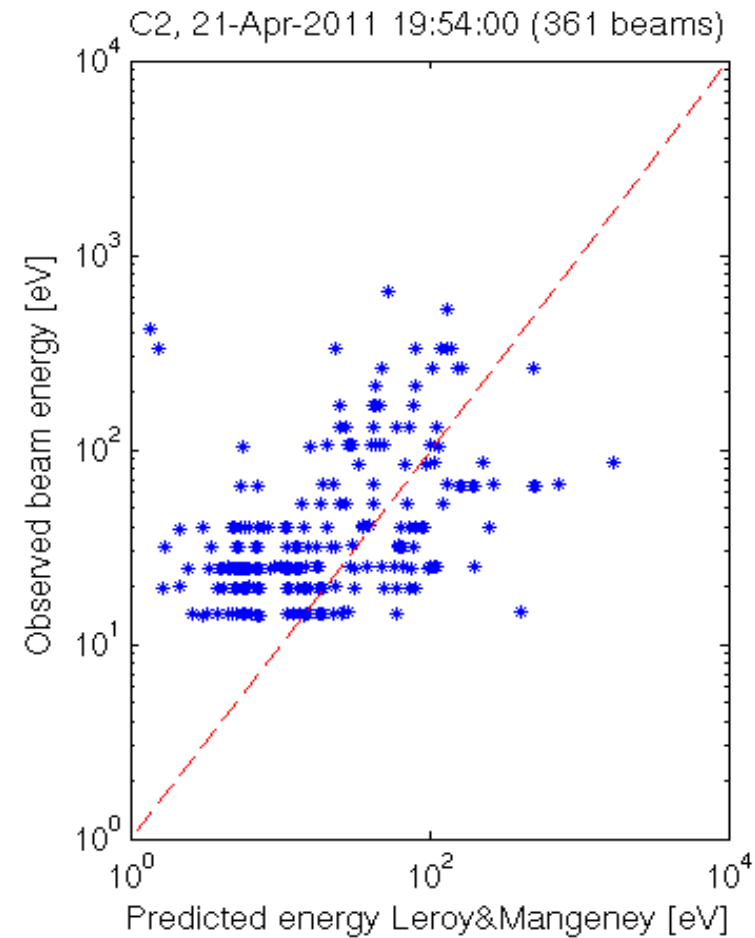
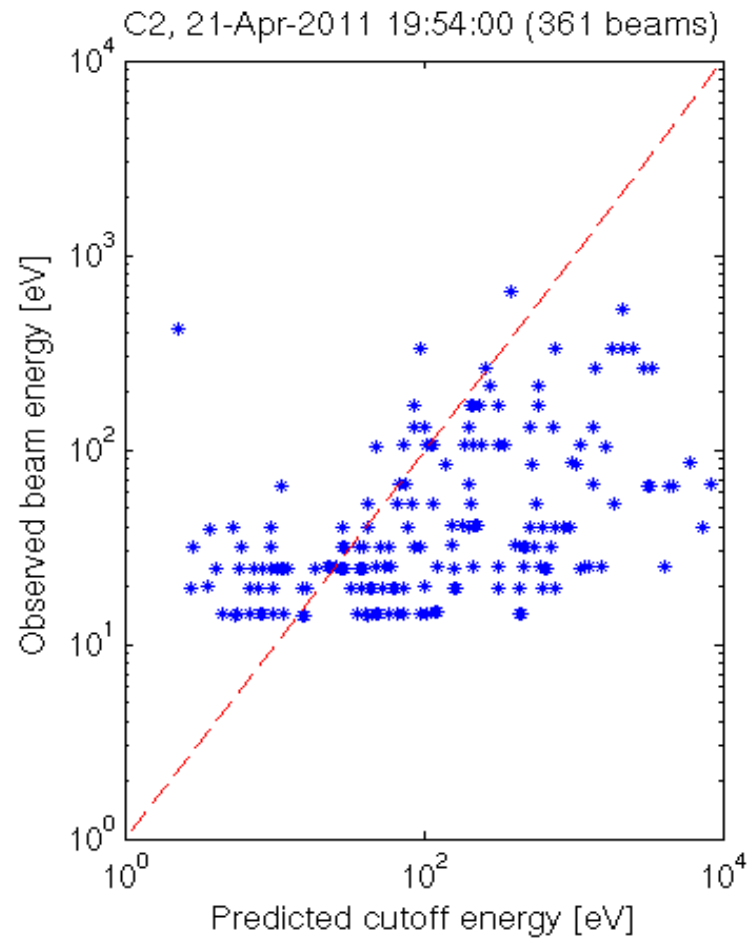
Beam velocity normalized to thermal velocity



Normalized to thermal velocity V_t

➔ $D_f < 2$ RE $\rightarrow V_b$ up to $10 V_t$
➔ $D_f > 2$ RE $\rightarrow V_b \sim 2-4 V_t$

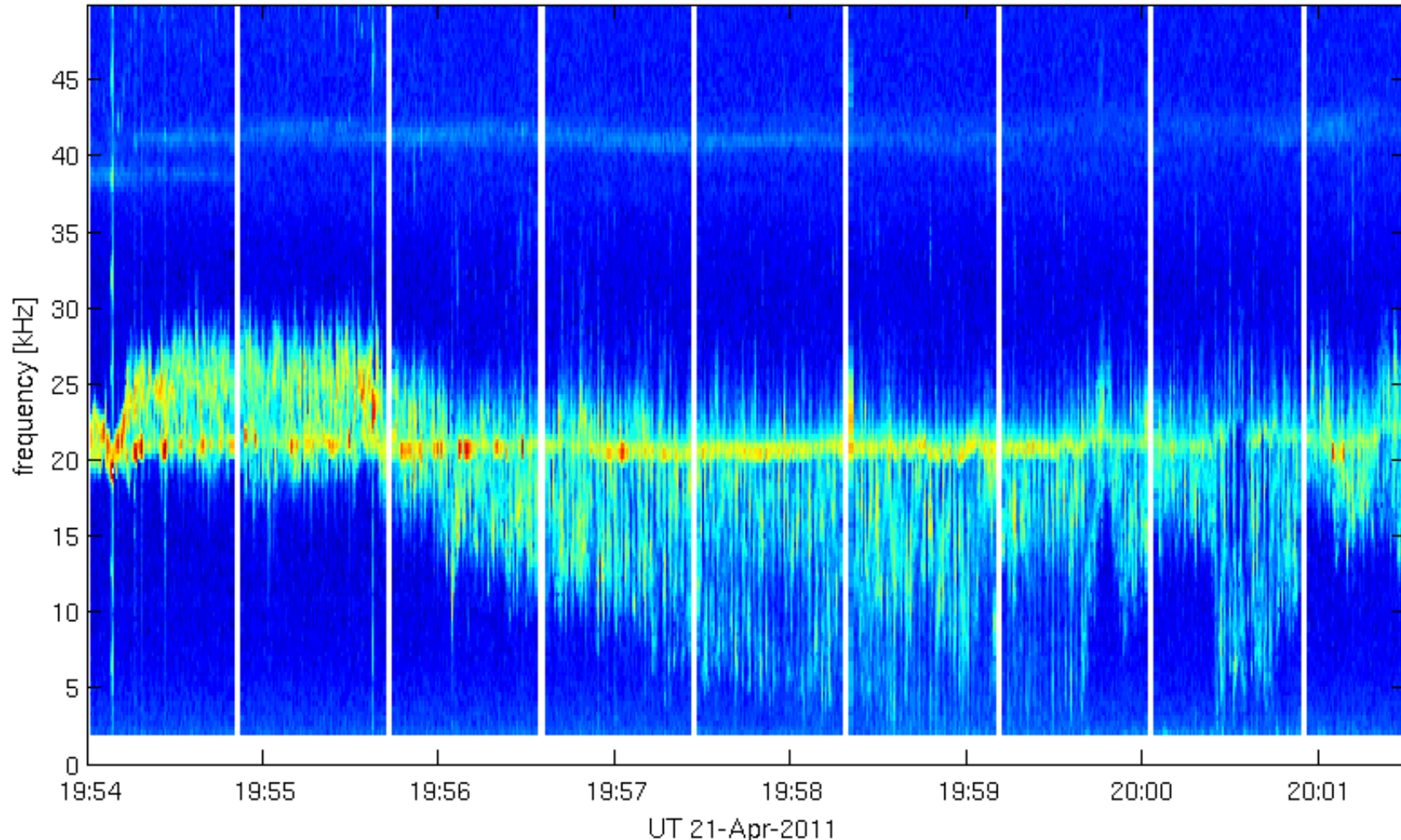
Comparison to models - statistics



- Comparison of measured and model beam speeds for moderate beam speeds.
 - ➔ Fitzenreiter et al., 1990 cutoff distribution model.
 - ➔ Leroy and Mangeny, 1984 energy.
- The measured and predicted value are comparable, but correlation is weak.

Foreshock wave structure

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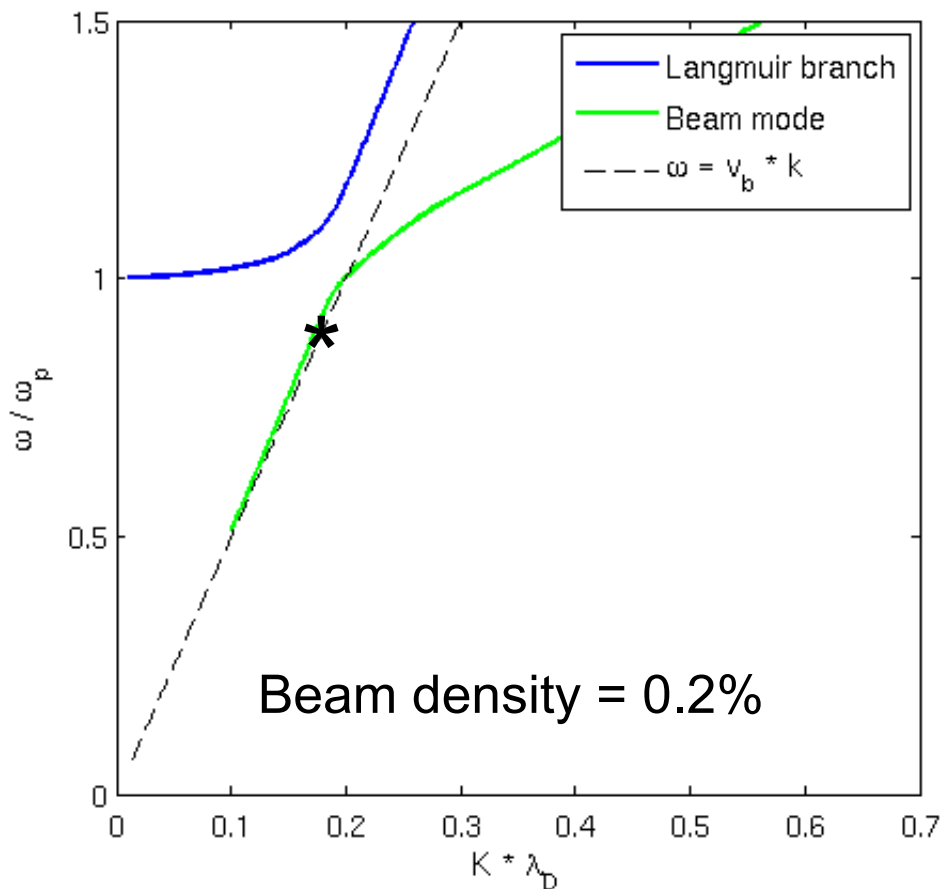


● Characteristic features of the foreshock spectrum:

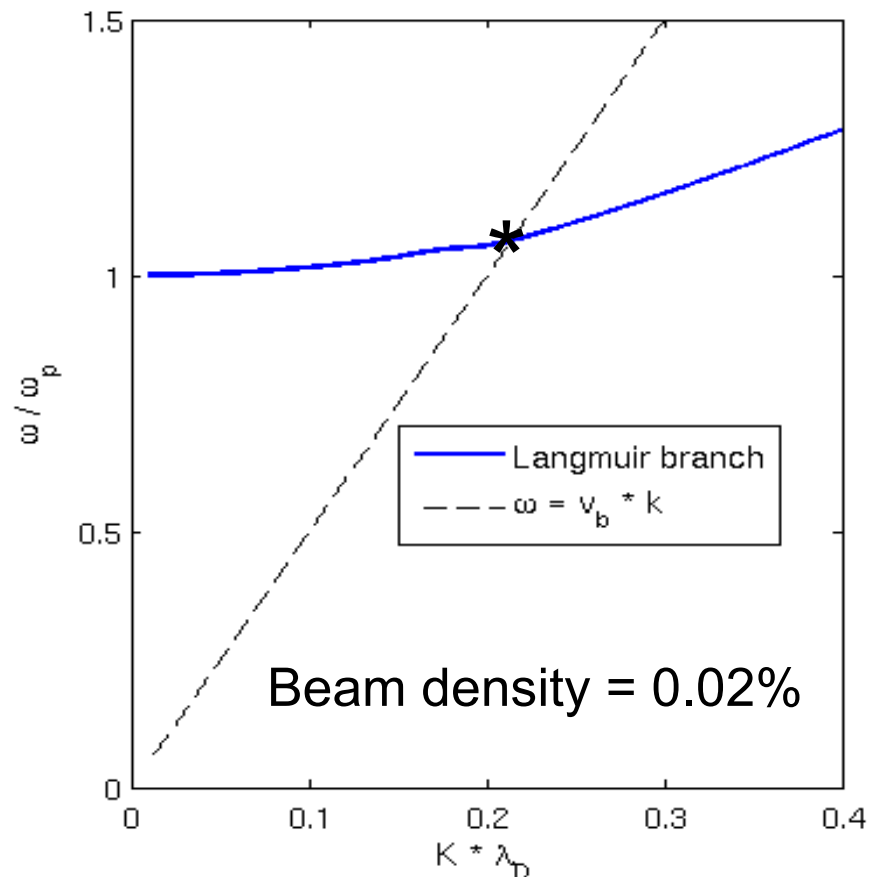
- ➔ Wave activity at plasma frequency.
- ➔ Waves far above and below the plasma frequency ($\Delta f/f$ up to 80%)

Plasma dispersion relation

$nb/n_0 = 0.0020, v_b/v_0 = 5.00, T_b/T_0 = 0.33$

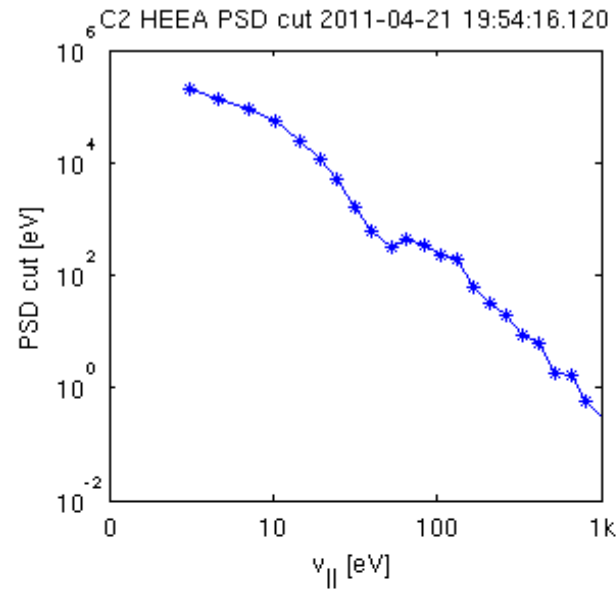
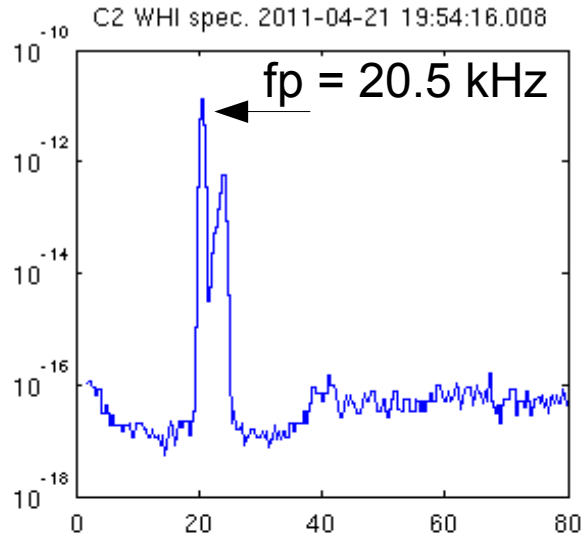


$nb/n_0 = 0.0002, v_b/v_0 = 5.00, T_b/T_0 = 0.33$



- Electrostatic dispersion relation for beam-plasma instability (Gary 1983)
 - ➔ Second solution – beam mode.
 - ➔ Langmuir wave branch modified
- Growth often on the beam mode
 - ➔ “Real” Langmuir waves only grow for fast and tenuous beams

Wave-particle analysis



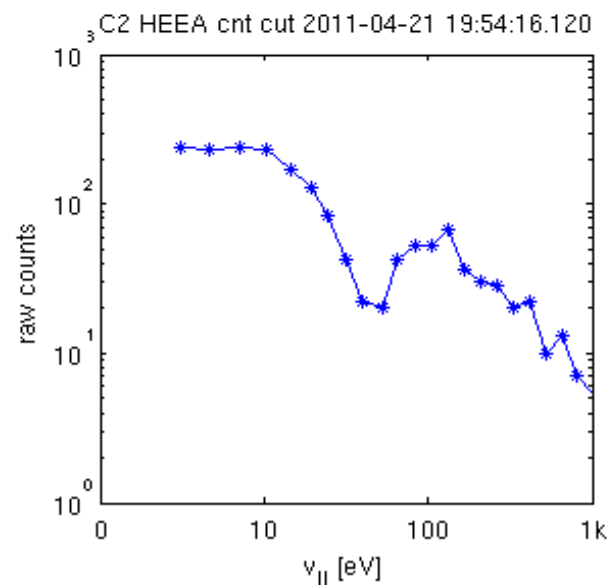
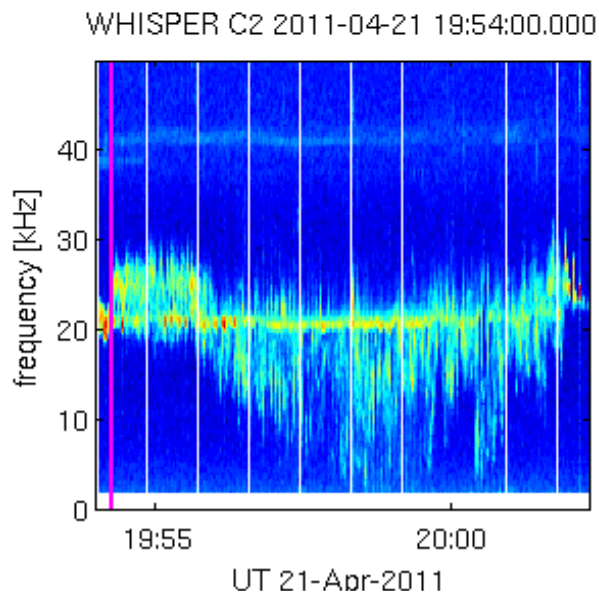
- Fitting of two Maxwellians employed to estimate beam properties

$$V_b = 86 \text{ eV} = 4.1 V_t$$

$$n_b/n_0 = 0.25\%$$

$$T_b/T_0 = 0.4$$

These 3 parameters are sufficient for linear dispersion analysis.

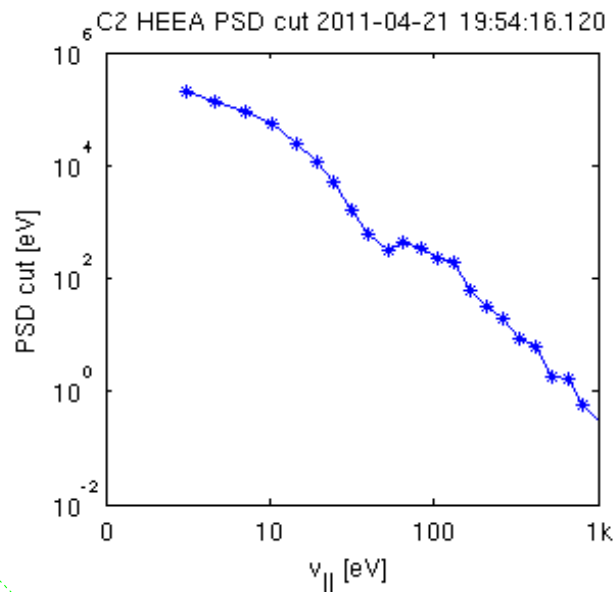
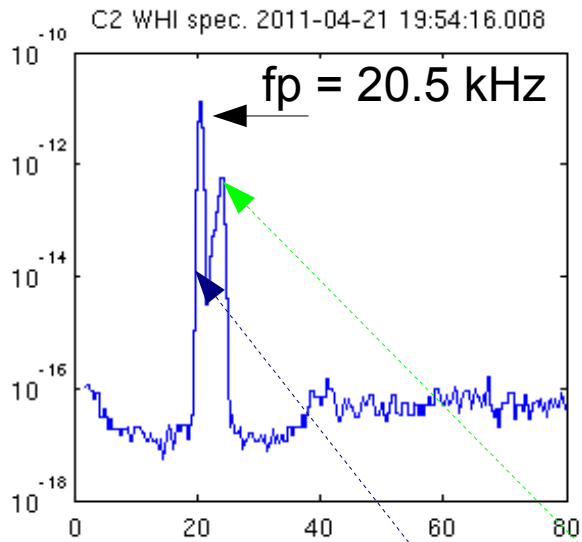


- Foreshock geometry:

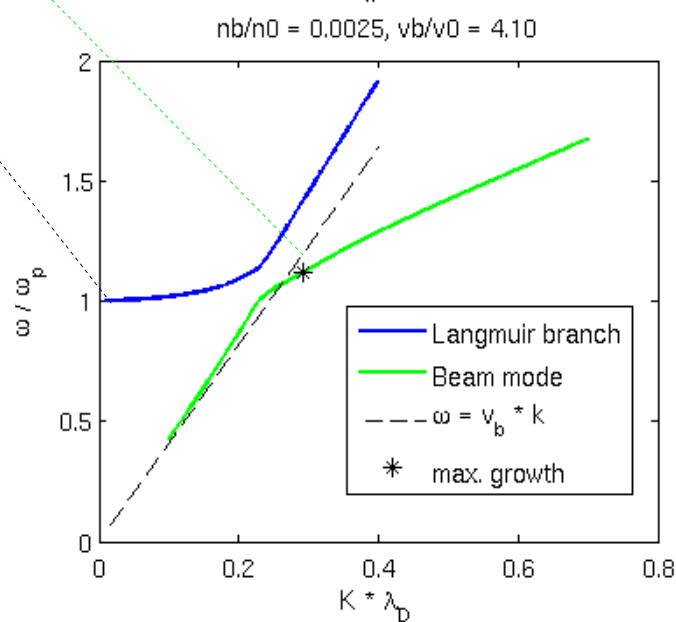
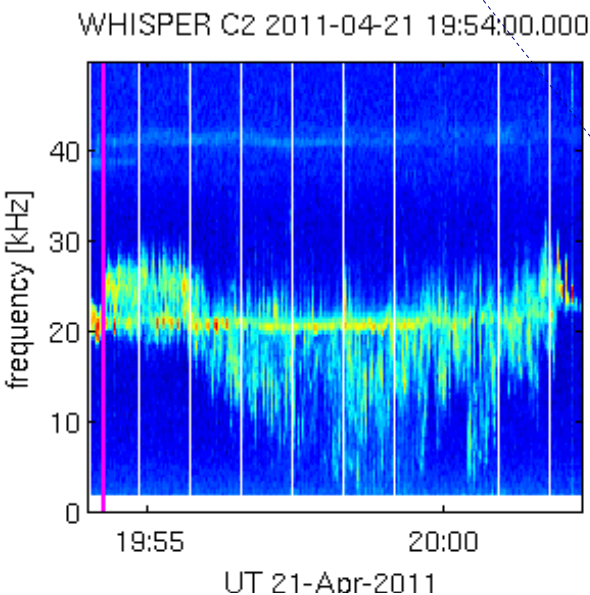
➔ The beam corresponds to $\Theta_{Bn} \sim 87^\circ$ source.

➔ Leroy & Mangeney model predicts $\sim 570 \text{ eV}$ (near 90° , model is very sensitive to Θ_{Bn} value).

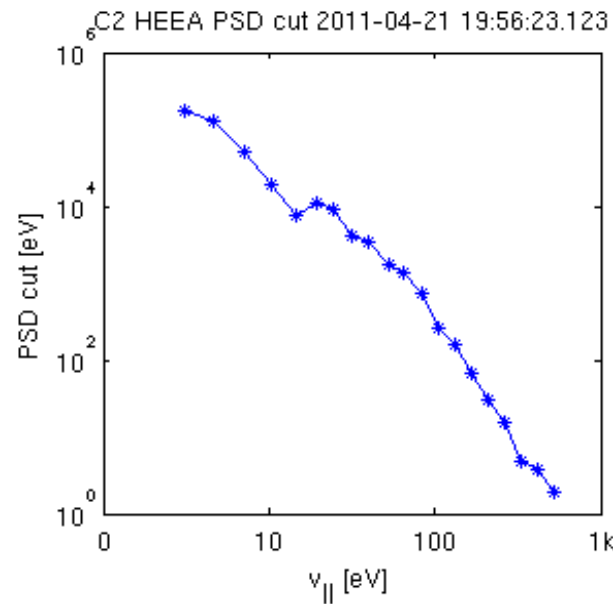
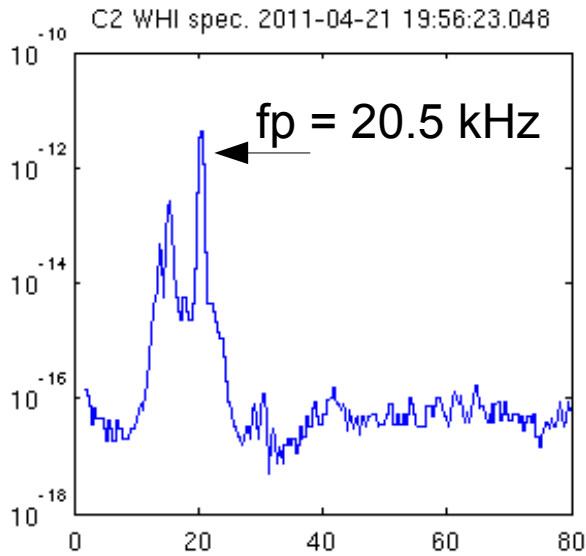
Wave-particle analysis



- Linear dispersion relation solved for measured parameters.
- Langmuir branch
 - ➔ No growth
 - ➔ Low damping at small k
- Beam mode
 - ➔ growth at $k \cdot \lambda_d \approx 0.3$
- Growth on the beam mode branch
 - ➔ Maximum calculated growth at 23.2 kHz
 - ➔ Peak observed at 24.3 kHz
- Oscillations near ω_p likely on Langmuir branch.



Wave-particle analysis



- Drifting Maxwellian fit:

$V_b = 31 \text{ eV} = 2.6 V_t$

$n_b/n_0 = 2.5\%$

$T_b/T_0 = 0.3$

- Foreshock geometry:

➔ The beam corresponds to $\Theta_{Bn} \sim 75^\circ$ source.

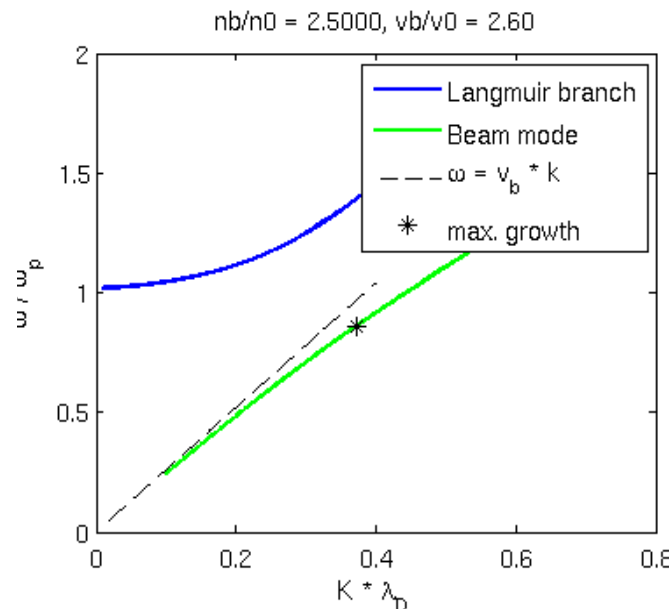
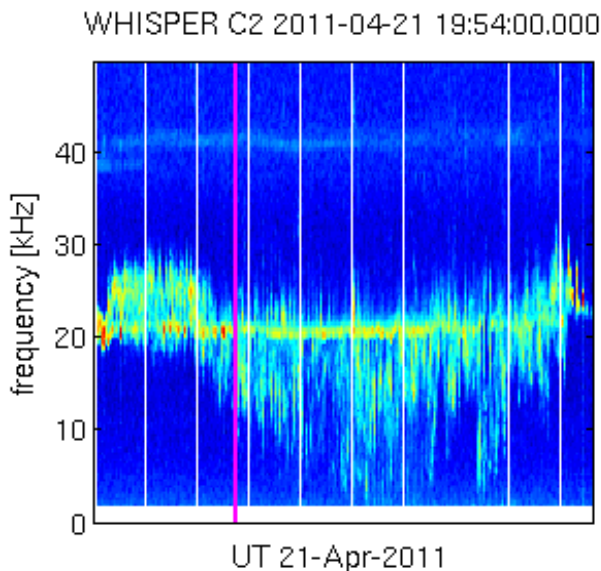
➔ Leroy & Mangeney model predicts $\sim 33 \text{ eV}$

- Growth below ω_p

➔ Max growth at 17.6 kHz

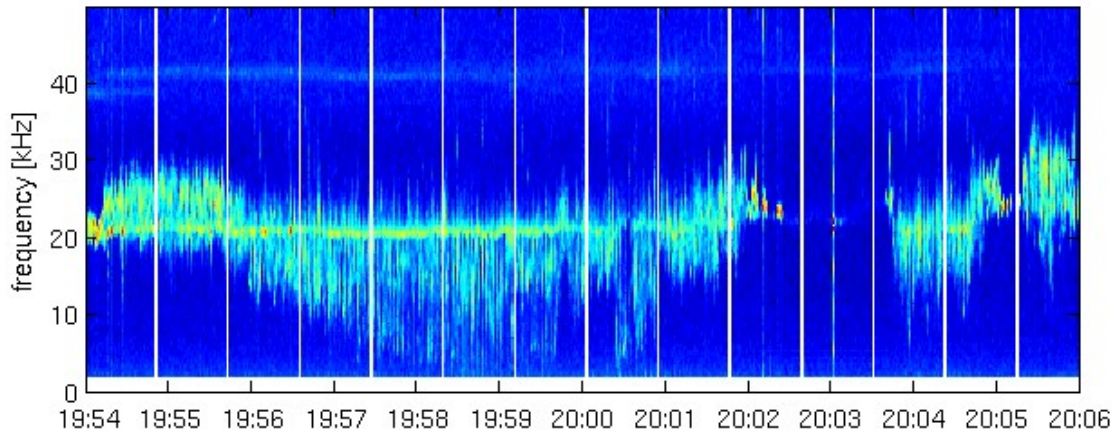
➔ Peak at 15.5 kHz

➔ “Rough” agreement only. Can probably be improved by a better distribution model.

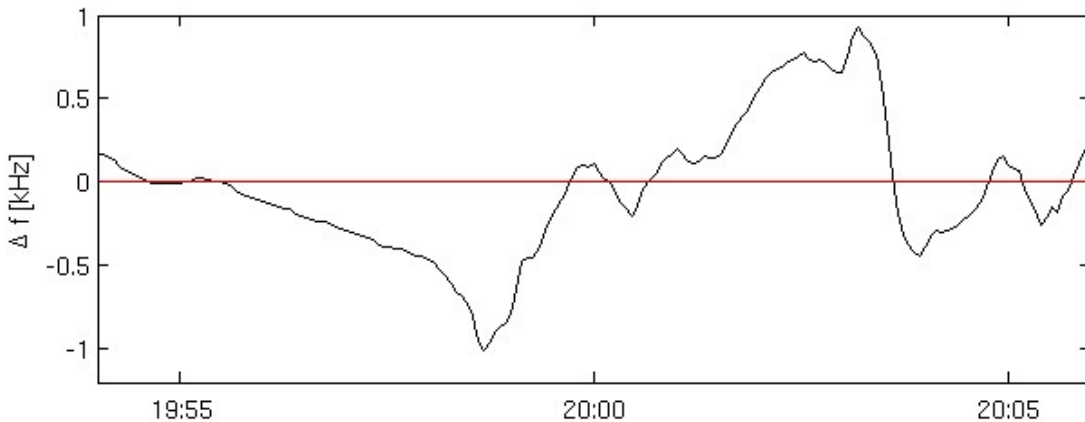


Role of the Doppler shift

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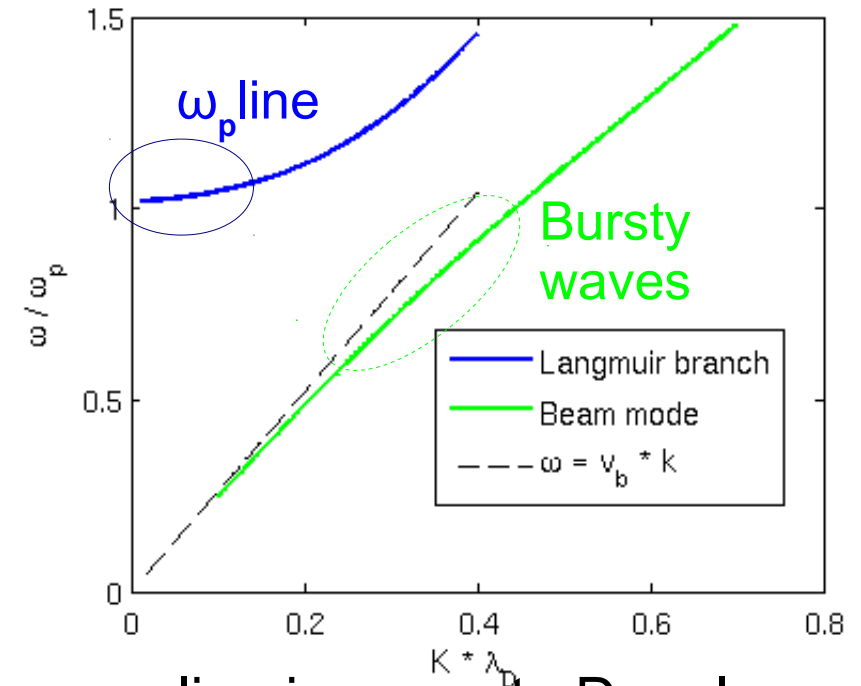


Predicted Doppler shift assuming $k \cdot \lambda_D = 0.4$



Predicted Doppler shift ($k \cdot \lambda_D$)

$nb/n0 = 0.0350, vb/v0 = 2.60$



- ω_p line immune to Doppler shift
 - ➔ Very low k waves
- For beam modes, Doppler shift up to 1 kHz
 - ➔ Direction consistent
 - ➔ Cannot account for all variation

Conclusions

- Systematic measurement of foreshock electron beam energies performed using PEACE data
 - ➔ Beams up to 1 keV observed consistently
 - ➔ Dependence on foreshock depth consistent with Leroy & Mangeney 1984 model.
- Broadband waves below ω_p correspond to slower beams [Lacombe et al., 1985, Cairns & Fung, 1988].
- Dispersion relation for electrostatic beam plasma instability solved for selected distributions.
 - ➔ Growth on the beam-mode branch at moderate k
 - ➔ Low k Langmuir oscillations at plasma frequency observed
 - ➔ Rough agreement between calculated and observed values
- Doppler shift significant, but cannot account for the variation in frequency.